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Әль-фараби атындағы Қазақ ұлттық университетінің

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## DIRICHLET PROBLEM IN A CYLINDRICAL AREA FOR ONE CLASS OF MULTIDIMENSIONAL ELLIPTIC-PARABOLIC EQUATIONS

**Abstract.** Boundary-value problems for degenerate elliptic-parabolic equations on the plane are studied quite well ([1]). The correctness of the Dirichlet problem for degenerate multidimensional elliptic-parabolic equations with degeneration of type and order was established in [3]. In the work for multidimensional elliptic-parabolic equations with degeneration of type and order, the solvability is shown and an explicit form of the classical solution of the Dirichlet problem is obtained.

**Keywords:** solvability, mixed problem, multidimensional elliptic-parabolic equations, Bessel function.

### Problem statement and result

Let  $\Omega_{\alpha\beta}$  – the cylindrical area of the Euclidean space of  $E_{m+1}$  points  $(x_1, \dots, x_m, t)$  bounded by a cylinder  $\Gamma = \{(x, t) : |x| = 1\}$ , planes  $t = \alpha > 0$  and  $t = \beta < 0$ , where  $|x|$  – is the length of a vector  $x = (x_1, \dots, x_m)$ .

Denote by  $\Omega_\alpha$  and  $\Omega_\beta$  parts  $\Omega_{\alpha\beta}$  – of the area and  $\Gamma_\alpha, \Gamma_\beta$  – through parts of the surface  $\Gamma$ , lying in the half-spaces  $t > 0$  and  $t < 0$ ,  $\sigma_\alpha$  – the upper and  $\sigma_\beta$  – lower base area  $\Omega_{\alpha\beta}$ .

Let  $S$  – further the common part of the borders of the areas  $\Omega_\alpha$  and  $\Omega_\beta$  representing the  $\{t = 0, 0 < |x| < 1\}$  set in  $E_m$ .

In the area  $\Omega_{\alpha\beta}$ , we consider degenerate multidimensional hyperbolic-parabolic equations

$$0 = \begin{cases} p_1(t)\Delta_x u - p_2(t)u_{tt} + \sum_{i=1}^m a_i(x, t)u_{x_i} + b(x, t)u_t + c(x, t)u = 0, t > 0, \\ g(t)\Delta_x u - u_t + \sum_{i=1}^m d_i(x, t)u_{x_i} + e(x, t)u, t < 0, \end{cases} \quad (1)$$

where  $p_i(t) > 0$  at  $t > 0$ ,  $p_i(0) = 0$ ,  $p_i(t) \in C([0, \alpha])$ ,  $g(t) > 0$  at  $t < 0$ , and may vanish when  $t = 0$ ,  $g(t) \in C[\beta, 0]$ , a  $\Delta_x$  – Laplace operator with variables  $x_1, \dots, x_m$ ,  $m \geq 2$ .

In the future, it is convenient for us to move from the Cartesian coordinates  $x_1, \dots, x_m, t$  to spherical  $r, \theta_1, \dots, \theta_{m-1}, t, r \geq 0, 0 \leq \theta_{m-1} < 2\pi, 0 \leq \theta_i \leq \pi, i = 1, 2, \dots, m-2, \theta = (\theta_1, \dots, \theta_{m-1})$ .

**Problem 1 (Dirichlet).** Find a solution to the equation (1) in the area of  $\Omega_{\alpha\beta}$  at  $t \neq 0$ , from the class  $C^1(\overline{\Omega_{\alpha\beta}}) \cap C^2(\Omega_\alpha \cup \Omega_\beta)$ , satisfying boundary conditions

$$u|_{\sigma_\alpha} = \varphi_1(r, \theta), \quad u|_{\Gamma_\alpha} = \psi_1(t, \theta), \quad (2)$$

$$u|_{\Gamma_\beta} = \psi_2(t, 0), \quad u|_{\sigma_\beta} = \varphi_2(t, \theta). \quad (3)$$

wherein  $\varphi_1(1, \theta) = \psi_1(\alpha, \theta), \psi_1(0, \theta) = \psi_2(0, \theta), \psi_2(\beta, \theta) = \varphi_2(1, \theta)$ .

Let  $\{Y_{n,m}^k(\theta)\}$  - system of linearly independent spherical functions of order  $n$ ,  $1 \leq k \leq k_n$ ,  $(m-2)!n!k_n = (n+m-3)!(2n+m-2), W_2^l(S), l = 0, 1, \dots$  - Sobolev space. Takes place ([4]).

**Lemma 1.** Let  $f(r, \theta) \in W_2^l(S)$ . If  $l \geq m-1$ , that row

$$f(r, \theta) = \sum_{n=0}^{\infty} \sum_{k=1}^{k_n} f_n^k(r) Y_{n,m}^k(\theta), \quad (4)$$

as well as series derived from it by order differentiation  $p \leq l - m + 1$ , converge absolutely and evenly.

**Lemma 2.** In order to  $f(r, \theta) \in W_2^l(S)$ , it is necessary and sufficient that the coefficients of the series (4) satisfy the inequalities.

$$|f_0^1(r)| \leq c_1, \quad \sum_{n=1}^{\infty} \sum_{k=1}^{k_n} n^{2l} |f_n^k(r)|^2 \leq c_2, \quad c_1, c_2 = const.$$

Through  $\tilde{d}_{in}^k(r, t), d_{in}^k(r, t), \tilde{e}_n^k(r, t), \tilde{d}_n^k(r, t), \rho_n^k, \bar{\varphi}_{1n}^k(r), \bar{\varphi}_{2n}^k(r), \psi_{1n}^k(t), \psi_{2n}^k(t)$ , denote the coefficients of the series (4), respectively functions  $d_i(r, \theta, t)\rho(\theta), d_i \frac{x_i}{r} \rho, e(r, \theta, t)\rho, d(r, \theta, t)\rho, \rho(\theta), i = 1, \dots, m, \varphi_1(r, \theta), \varphi_2(r, \theta), \psi_1(t, \theta), \psi_2(t, \theta)$ , and  $\rho(\theta) \in C^\infty(H)$ ,  $H$ -unit sphere in  $E_m$ .

$$\text{Let } \frac{a_i(r, \theta, t)}{g_2(t)}, \frac{b(r, \theta, t)}{g_2(t)}, \frac{c(r, \theta, t)}{g_2(t)} \in W_2^l(\Omega_\alpha) \subset C(\bar{\Omega}_\alpha), d_i(r, \theta, t),$$

$$c(r, \theta, t) \in W_2^l(\Omega_\beta), i = 1, \dots, m, l \geq m+1, c(r, \theta, t) \leq 0, \forall (r, \theta, t) \in \Omega_\alpha, e(r, \theta, t) \in \Omega_\beta.$$

Then fair  
Theorem.

If  $\varphi_1(r, \theta), \varphi_2(r, \theta) \in W_2^l(S), \psi_1(t, \theta) \in W_2^p(\Gamma_\alpha), \psi_2(t, \theta) \in W_2^l(\Gamma_\beta), l > \frac{3m}{2}$ , then problem 1 is solvable.

**Proof of the theorem.** First, let us rock the solvability of problem (1), (3). In spherical coordinates of equation (1) in the area  $\Omega_\beta$  has the appearance

$$Lu \equiv g(t)(u_{rr} + \frac{m-1}{r}u_r - \frac{1}{r^2}\delta u) - u_{tt} + \sum_{i=1}^m d_i(r, \theta, t)u_{x_i} + e(r, \theta, t)u = 0, \quad (5)$$

$$\delta \equiv -\sum_{j=1}^{m-1} \frac{1}{g_j \sin^{m-j-1} \theta_j} \frac{\partial}{\partial \theta_j} (\sin^{m-j-1} \theta_j \frac{\partial}{\partial \theta_j}), g_1 = 1, g_j = (\sin \theta_1 \dots \sin \theta_{j-1})^2, j > 1.$$

It is known [4] that the spectrum of the operator  $\delta$  consists of own numbers  $\lambda_n = n(n+m-2), n = 0, 1, \dots$  each of which corresponds  $k_n$  orthonormal functions  $Y_{n,m}^k(\theta)$ .

The desired solution to problem 1 in the field  $\Omega_\beta$  we will look in the form

$$u(r, \theta, t) = \sum_{n=0}^{\infty} \sum_{k=1}^{k_n} \bar{u}_n^k(r, t) Y_{n,m}^k(\theta), \tag{6}$$

where  $\bar{u}_n^k(r, t)$  - functions to be defined.

Substituting (6) в (5), then multiplying the resulting expression by  $\rho(\theta) \neq 0$ , and integrating over a single sphere H, for  $\bar{u}_n^k$  will get [5-7]

$$\begin{aligned} &g(t) \rho_0^1 \bar{u}_{0rr}^1 + \rho_0^1 \bar{u}_{0tt}^1 + \left(\frac{m-1}{r} g(t) \rho_0^1 + \sum_{i=1}^m d_{i0}^1\right) u_{0r}^1 + \\ &+ \sum_{n=1}^{\infty} \sum_{k=1}^{k_n} \left\{ g(t) \rho_n^k \bar{u}_{nrr}^k + \rho_n^k \bar{u}_{ntt}^k + \left(\frac{m-1}{r} g(t) \rho_n^k + \sum_{i=1}^m d_{in}^k\right) \bar{u}_{nr}^k \right. \\ &\left. + [\tilde{c}_n^k - \lambda_n \frac{\rho_n^k}{r^2} g(t) + \sum_{i=1}^m (\tilde{d}_{in-1}^k - n d_{in}^k)] \bar{u}_n^k \right\} = 0. \end{aligned} \tag{7}$$

Now consider the infinite system of differential equations

$$g(t) \rho_0^1 \bar{u}_{0rr}^1 + \rho_0^1 u_{0tt}^1 + \frac{m-1}{r} g(t) \rho_0^1 \bar{u}_{0r}^1 = 0, \tag{8}$$

$$\begin{aligned} &g(t) \rho_1^k \bar{u}_{1rr}^k - \rho_1^k \bar{u}_{1t}^k + \frac{m-1}{r} g(t) \rho_1^k \bar{u}_{1r}^k - \frac{\lambda_1}{r^2} g(t) \rho_1^k \bar{u}_1^k = \\ &= -\frac{1}{k_1} \left( \sum_{i=1}^m d_{i0}^1 \bar{u}_{0r}^1 + \tilde{e}_0^1 \bar{u}_o^1 \right), \quad n = 1, \quad k = \overline{1, k_1}, \end{aligned} \tag{9}$$

$$\begin{aligned} &g(t) \rho_n^k \bar{u}_{nrr}^k - \rho_n^k \bar{u}_{nt}^k + \frac{m-1}{r} g(t) \rho_n^k \bar{u}_{nr}^k - \frac{\lambda_n}{r^2} g(t) \rho_n^k \bar{u}_n^k = \\ &= -\frac{1}{k_n} \sum_{k=1}^{k_{n-1}} \left\{ \sum_{i=1}^m d_{in-1}^k \bar{u}_{n-1r}^k + [\tilde{e}_{n-1}^k + \sum_{i=1}^m (\tilde{d}_{in-2}^k - (n-1) d_{in-1}^k)] \bar{u}_{n-1}^k \right\}, \\ &k = \overline{1, k_n}, \quad n = 2, 3, \dots \end{aligned} \tag{10}$$

Summing up the equation (8) from 1 before  $k_1$ , and the equation (9)- from 1 before  $k_n$ , and then adding the resulting expressions together with (7), come to the equation (6).

It follows that if  $\{\bar{u}_n^k\}, k = \overline{1, k_n}, n = 0, 1, \dots$  system solution (7)-(9), then it is a solution to the equation (6).

It is easy to see that each equation of system (7) - (9) can be represented as

$$g(t) \left( \bar{u}_{nrr}^k + \frac{m-1}{r} \bar{u}_{nr}^k - \frac{\lambda_n}{r^2} \bar{u}_n^k \right) - u_{nt}^k = \bar{f}_n^k(r, t), \tag{11}$$

where  $\bar{f}_n^k(r, t)$  are determined from the previous equations of this system, while  $\bar{f}_0^1(r, t) \equiv 0$ .

Further, from the boundary condition (3), by virtue of (6), we will have

$$\bar{u}_n^k(r, \beta) = \bar{\varphi}_{2n}^k(r), \quad \bar{u}_n^k(1, t) = \psi_{2n}^k(t), \quad k = \overline{1, k_n}, \quad n = 0, 1, \dots \tag{12}$$

In (11), (12) replacing  $\bar{\mathcal{G}}_n^k(r, t) = \bar{u}_n^k(r, t) - \psi_{2n}^k(t)$ , will get

$$g(t) \left( \bar{\mathcal{G}}_{nrr}^k + \frac{m-1}{r} \bar{\mathcal{G}}_{nr}^k - \frac{\lambda_n}{r^2} \bar{\mathcal{G}}_n^k \right) - \bar{\mathcal{G}}_{nt}^k = \bar{f}_n^k(r, t), \tag{13}$$

$$\bar{\mathcal{G}}_n^k(r, \beta) = \varphi_{2n}^k(r), \quad \bar{\mathcal{G}}_n^k(1, t) = 0, \quad k = \overline{1, k_n}, \quad n = 0, 1, \dots \quad (14)$$

$$f_n^k(r, t) = \bar{f}_n^k(r, t) + \psi_{2nt}^k + \frac{\lambda_n g(t)}{r^2} \psi_{2n}^k, \quad \varphi_{2n}^k(r) = \bar{\varphi}_{2n}^k(r) - \psi_{2n}^k(\beta).$$

Replacing the variable  $\bar{\mathcal{G}}_n^k(r, t) = r^{\frac{(1-m)}{2}} \mathcal{G}_n^k(r, t)$  задачи (13), (14) we will lead to the following problem

$$L\mathcal{G}_n^k = g(t)(\mathcal{G}_{nrr}^k + \frac{\bar{\lambda}_n}{r^2} \mathcal{G}_n^k) - \mathcal{G}_{nt}^k = \tilde{f}_n^k(r, t), \quad (15)$$

$$\mathcal{G}_n^k(r, \beta) = \bar{\varphi}_{2n}^k(r), \quad \mathcal{G}_n^k(1, t) = 0, \quad \mathcal{G}_n^k(1, t) = 0, \quad (16)$$

$$\bar{\lambda}_n = \frac{[(m-1)(3-m) - 4\lambda_n]}{4}, \quad \tilde{f}_n^k(r, t) = r^{\frac{(m-1)}{2}} f_n^k(r, t),$$

$$\tilde{\varphi}_{2n}^k(r) = r^{\frac{(m-1)}{2}} \varphi_{2n}^k(r).$$

The solution of the problem (15), (16) is sought in the form

$$\mathcal{G}_n^k(r, t) = \mathcal{G}_{1n}^k(r, t) + \mathcal{G}_{2n}^k(r, t), \quad (17)$$

where  $\mathcal{G}_{1n}^k(r, t)$  the solution of the problem

$$L\mathcal{G}_{1n}^k = \tilde{f}_n^k(r, t), \quad (18)$$

$$\mathcal{G}_{1n}^k(r, \beta) = 0, \quad \mathcal{G}_{1n}^k(1, t) = 0, \quad (19)$$

where  $\mathcal{G}_{2n}^k(r, t)$  the solution of the problem

$$L\mathcal{G}_{1n}^k = 0, \quad (20)$$

$$\mathcal{G}_{2n}^k(r, \beta) = \tilde{\varphi}_{2n}^k(r), \quad \mathcal{G}_{2n}^k(1, t) = 0, \quad (21)$$

The solution to the above problems, we consider in the form

$$\mathcal{G}_n^k(r, t) = \sum_{s=1}^{\infty} R_s(r) T_s(t), \quad (22)$$

at the same time let

$$\tilde{f}_n^k(r, t) = \sum_{s=1}^{\infty} a_{ns}^k(t) R_s(r), \quad \tilde{\varphi}_{2n}^k(r) = \sum_{s=1}^{\infty} b_{ns}^k R_s(r). \quad (23)$$

Substituting (22) into (18), (19), taking into account (23), we obtain

$$R_{srr} + \frac{\bar{\lambda}_n}{r^2} R_s + \mu_{s,n} R_s = 0, \quad 0 < r < 1, \quad (24)$$

$$R_s(1) = 0, \quad |R_s(0)| < \infty, \quad (25)$$

$$T_{st} - \mu_{s,n} g(t) T_s(t) = -a_{ns}^k(t), \quad \beta < t < 0, \quad (26)$$

$$T_s(\beta) = 0. \quad (27)$$

A limited solution to problem (24), (25) is ([8])



$$R_s(r) = \sqrt{r} J_\nu(\mu_{s,n} r), \tag{28}$$

where  $\nu = \frac{n + (m - 2)}{2}$ ,  $\mu_{s,n}$  - zeros of the Bessel function of the first kind  $J_\nu(z)$ ,  $\mu = \mu_{s,n}^2$ .

The solution to problem (26), (27) is

$$T_{s,n}(t) = (\exp(-\mu_{s,n}^2 \int_0^t g(\xi) d\xi)) \int_t^\beta g(\xi) (\exp \mu_{s,n}^2 \int_0^\xi g(\xi_1) d\xi_1) d\xi. \tag{29}$$

Substituting (28) into (23) we get

$$r^{\frac{1}{2}} \tilde{f}_n^k(r, t) = \sum_{s=1}^\infty a_{ns}^k(t) J_\nu(\mu_{s,n} r), \quad r^{\frac{1}{2}} \tilde{\varphi}_{1n}^k(r) = \sum_{s=1}^\infty b_{ns}^k J_\nu(\mu_{s,n} r), \quad 0 < r < 1. \tag{30}$$

Rows (30) - Fourier-Bessel series expansions ([9]), if

$$a_{ns}^k(t) = 2 [J_{\nu+1}(\mu_{s,n})]^{-2} \int_0^1 \sqrt{\xi} \tilde{f}_n^k(\xi, t) J_\nu(\mu_{s,n} \xi) d\xi. \tag{31}$$

$$b_{ns}^k = 2 [J_{\nu+1}(\mu_{s,n})]^{-2} \int_0^1 \sqrt{\xi} \tilde{\varphi}_{2n}^k(\xi) J_\nu(\mu_{s,n} \xi) d\xi, \tag{32}$$

where  $\mu_{s,n}$   $s = 1, 2, \dots$  - positive zeros of the Bessel function  $J_\nu(z)$ , located in ascending order of magnitude.

Of (22), (28), (29) get the solution to the problem (18), (19)

$$\mathcal{G}_{1n}^k(r, t) = \sum_{s=1}^\infty \sqrt{r} T_{s,n}(t) J_\nu(\mu_{s,n} r), \tag{33}$$

where  $a_{ns}^k(t)$  determined from (31).

Next, substituting (22) в (20), (21), taking into account (23), will have

$$T_{st} - \mu_{s,n}^2 g(t) T_s = 0, \quad \beta < t < 0, \quad T_s(\beta) = b_{ns}^k,$$

which solution is

$$T_{s,n}(t) = b_{ns}^k \exp(\mu_{s,n}^2 \int_t^\beta g(\xi) d\xi). \tag{34}$$

From (28), (34) we get

$$\mathcal{G}_{2n}^k(r, t) = \sum_{s=1}^\infty b_{ns}^k \sqrt{r} (\exp \mu_{s,n}^2 \int_t^\beta g(\xi) d\xi) J_\nu(\mu_{s,n} r), \tag{35}$$

where  $b_{ns}^k$  are from (32).

Therefore, first solving the problem (8), (12) (n=0), and then (9), (12) (n=1) etc. let's find everything  $\mathcal{G}_n^k(r, t)$  of (17), where  $\mathcal{G}_{2n}^k(r, t)$  are determined from (33), (35).

So, in the area  $D_\alpha$ , takes place

$$\int_H \rho(\theta) L_1 u dH = 0. \tag{36}$$

Let  $f(r, \theta, t) = R(r) \rho(\theta) T(t)$ , and  $R(r) \in V_0, V_0$  - tight in  $L_2((0, 1))$ ,  $\rho(\theta) \in C^\infty(H)$  - tight in  $L_2(H)$ ,  $T(t) \in V_1, V_1$  - tight in  $L_2((\beta, 0))$ . Then  $f(r, \theta, t) \in V, V = V_0 \otimes H \otimes V_1$  - tight in  $L_2(\Omega_\beta)$  [10].

From here and from (36), it follows that

$$\int_{\Omega_\beta} f(r, \theta, t) L_1 u d\Omega_\beta = 0$$

and

$$L_1 u = 0, \quad \forall (r, \theta, t) \in \Omega_\beta.$$

Thus, by solving the problem (1), (3) in the field  $\Omega_\beta$  is the function

$$u(r, \theta, t) = \sum_{n=0}^{\infty} \sum_{k=1}^{k_n} \{ \psi_{2n}^k(t) + r^{\frac{(1-m)}{2}} [\mathcal{G}_{1n}^k(r, t) + \mathcal{G}_{2n}^k(r, t)] \} Y_{n,m}^k(\theta), \quad (37)$$

where  $\mathcal{G}_{1n}^k(r, t), \mathcal{G}_{2n}^k(r, t)$  are from (33), (35).

Given the formula ([9]):

$$2J'_\nu(z) = J_{\nu-1}(z) - J_{\nu+1}(z), \quad \text{ratings [11,4]}$$

$$J_\nu(z) = \sqrt{\frac{2}{\pi z}} \cos\left(z - \frac{\pi}{2}\nu - \frac{\pi}{4}\right) + O\left(\frac{1}{z^{\frac{3}{2}}}\right), \quad \nu \geq 0,$$

$$|k_n| \leq c_1 n^{m-2}, \quad \left| \frac{\partial^l}{\partial \theta_j^l} Y_{n,m}^k(\theta) \right| \leq c_2 n^{\frac{m}{2}-1+l}, \quad j = \overline{1, m-1}, l = 0, 1, \dots, \quad (38)$$

as well as lemmas, restrictions on the coefficients of equation (1) and on given functions  $\varphi_1(r, \theta), \varphi_2(r, \theta), \psi_1(t, \theta), \psi_2(t, \theta)$  can be shown that received solution (37) belongs to the class  $C^1(\overline{\Omega}_\beta) \cap C^2(\Omega_\beta)$ .

$$u(r, \theta, 0) = \tau(r, \theta) = \sum_{n=0}^{\infty} \sum_{k=1}^{k_n} \tau_n^k(r) Y_{n,m}^k(\theta), \quad (39)$$

$$\tau_n^k(r) = \psi_{2n}^k(0) + \sum_{s=1}^{\infty} r^{\frac{(2-m)}{2}} \left[ \int_0^\beta a_{ns}^k(\xi) (\exp \mu_{s,n}^2 \int_0^\xi g(\xi_1) d\xi_1) d\xi + b_{ns}^k \exp(\mu_{s,n}^2 \int_0^\beta g(\xi) d\xi) J_\nu(\mu_{s,n} r) \right].$$

From (30) - (33), (35), and also from the lemmas, it follows that  $\tau(r, \theta) \in W_2^l(S)$ ,  $l > \frac{3m}{2}$ .

Thus, taking into account the boundary conditions (2) and (39), we arrive at  $\Omega_\beta$  to the Dirichlet problem for an elliptic equation.

$$L_2 u \equiv p_1(t) \Delta_x u + p_2(t) u_{tt} + \sum_{i=1}^m a_i(r, \theta, t) u_{x_i} + b(r, \theta, t) u_t + c(r, \theta, t) u = 0, \quad (40)$$

with data

$$u|_{S_0} = \tau(r, \theta), \quad u|_{\Gamma_\alpha} = \psi_1(t, \theta), \quad u|_{\sigma_\alpha} = \varphi_1(r, \theta), \quad (41)$$

having a solution ([12]).

Hence the solvability of the problem 1 is established.

**The uniqueness of the solution to problem 1.** First, we consider problem (1), (3) in the area  $\Omega_\beta$  and prove its uniqueness of the solution. To do this, we first construct the solution of the first boundary value problem for the equation

$$L_1^* \mathcal{G} \equiv g(t) \Delta_x \mathcal{G} + \mathcal{G}_t - \sum_{i=1}^m d_i \mathcal{G}_{x_i} + d \mathcal{G} = 0, \tag{5^*}$$

with data

$$\mathcal{G}|_s = \tau(r, \theta) \sum_{n=0}^{\infty} \sum_{k=1}^{k_n} \bar{\tau}_n^k(r) Y_{n,m}^k(\theta), \quad \mathcal{G}|_{\Gamma_\beta} = 0, \tag{42}$$

where  $d(x, t) = e - \sum_{i=1}^m d_{ix_i}$ ,  $\bar{\tau}_n^k(r) \in G, G$  - many functions  $\tau(r)$  from the class  $C([0, 1]) \cap C^1(0, 1)$ . Lots of  $G$  tight everywhere in  $L_2((0, 1))$  [10]. The solution to the problem (5\*), (42) will be sought in the form (6), where the functions  $\mathcal{G}_n^k(r, t)$  will be defined below. Then, similarly to item 2. functions  $\bar{\mathcal{G}}_n^k(r, t)$  satisfy the system of equations (8)-(10), where  $\bar{d}_{in}^k, d_{in}^k$  replaced respectively by  $-\bar{d}_{in}^k, -d_{in}^k$ , a  $\bar{\tau}_n^k$  on  $\tau_n^k$ ,  $i = 1, \dots, m, k = \overline{1, k_n}, n = 0, 1, \dots$ .

Further, from the boundary condition (42), by virtue of (6), we obtain

$$\bar{\mathcal{G}}_n^k(r, \theta) = \bar{\tau}_n^k(r), \quad \bar{\mathcal{G}}_n^k(1, t) = 0, \quad k = \overline{1, k_n}, \quad n = 0, 1, \dots \tag{43}$$

As previously noted, each equation of system (8) - (10) is represented as (11). Problem (11), (43) we will result in the following problem.

$$L \mathcal{G}_n^k = g(t) (\mathcal{G}_{nrr}^k + \frac{\bar{\lambda}_n}{r^2} \mathcal{G}_n^k) + \mathcal{G}_{nt}^k = \tilde{f}_n^k(r, t), \tag{15'}$$

$$\mathcal{G}_n^k(r, 0) = \tau_n^k(r), \quad \mathcal{G}_n^k(1, t) = 0, \tag{44}$$

$$\mathcal{G}_n^k(r) = r^{\frac{(m-1)}{2}} \bar{\mathcal{G}}_n^k(r, t), \quad \tilde{f}_n^k(r, t) = r^{\frac{(m-1)}{2}} \bar{f}_n^k(r, t), \quad \tau_n^k(r) = r^{\frac{(m-1)}{2}} \bar{\tau}_n^k(r).$$

The solution to problem (15), (44) will be sought in the form (17), where  $\mathcal{G}_n^k(1, t)$  - solution of the problem for equation (18) with the data  $\mathcal{G}_{2n}^k(r, t)$

$$\mathcal{G}_{1n}^k(r, 0) = 0, \quad \mathcal{G}_{1n}^k(1, t) = 0, \tag{45}$$

a -  $\mathcal{G}_{2n}^k(r, t)$  solution of the problem for equation (20) with the condition

$$\mathcal{G}_{2n}^k(r, 0) = 0, \quad \mathcal{G}_{2n}^k(1, t) = 0, \tag{46}$$

The solution of problems (18), (45) and (20), (46) respectively I have the form

$$\mathcal{G}_{1n}^k(r, t) = \sum_{s=1}^{\infty} \sqrt{r} (\exp(+\mu_{s,n}^2 \int_0^t g(\xi) d\xi)) (\int_0^t a_{ns}^k(\xi) (\exp(\mu_{s,n}^2 \int_0^\xi g(\xi_1) d\xi_1) J_\nu(\mu_{s,n} r)),$$

where

$$\tau_{s,n} = 2 [J_{\nu+1}(\mu_{s,n})]^{-2} \int_0^1 \sqrt{\xi} \tau_n^k(\xi) J_\nu(\mu_{s,n} \xi) d\xi, \quad \nu = \frac{n + (m - 2)}{2}.$$

Thus, solving the problem (5\*), (42) in the form of a series

$$u(r, \theta, t) = \sum_{n=0}^{\infty} \sum_{k=1}^{k_n} r^{\frac{(1-m)}{2}} [\mathcal{G}_{1n}^k(r, t) + \mathcal{G}_{2n}^k(r, t)] Y_{n,m}^k(\theta),$$

built, which by virtue of estimates (38) belongs to the class  $C(\bar{\Omega}_\beta) \cap C^2(\Omega_\beta)$ .

As a result of integration by area  $\Omega_\beta$  identity [13]

$$\mathcal{G}L_1 u - uL_1^* \mathcal{G} = -\mathcal{G}P(u) + uP(\mathcal{G}) - u\mathcal{G}Q,$$

where

$$P(u) = g(t) \sum_{i=1}^m u_{x_i} \cos(N^\perp, x_i), \quad Q = \cos(N^\perp, t) - \sum_{i=1}^m d_i \cos(N^\perp, x_i),$$

but  $N^\perp$  - internal normal to the border  $\partial\Omega_\beta$ , according to the Green formula we get

$$\int_S \tau(r, \theta) u(r, \theta, 0) ds = 0. \quad (47)$$

Since the linear span of a system of functions  $\{\bar{\tau}_n^k(r) Y_{n,m}^k(\theta)\}$  tight  $L_2(S)$  ([10]), then from (47) we conclude that  $u(r, \theta, 0) = 0, \forall (r, \theta) \in S$ . So on the principle of extremum for a parabolic equation (5) [14]  $u \equiv 0$  в  $\bar{\Omega}_\beta$ .

Next, from the Hopf principle ([15])  $u \equiv 0$  в  $\bar{\Omega}_\beta$ .

The theorem is proven completely.

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#### КӨП-ӨЛШЕМДІ ЭЛЛИПТИКО-ПАРАБОЛАЛЫҚ ТЕНДЕУЛЕРІНІҢ БІР КЛАСЫ БОЙЫНША ЦИЛИНДРЛІК ОБЛЫСЫНДА ДИРИХЛЕ ЕСЕБІ

**Аннотация.** Жазықтықтағы эллиптико-параболикалық тендеулер үшін шеттік есептер өте жақсы зерттелген ([1]). Дирихле есебінің корректілігі түрі мен ретті алып тұратын көп өлшемді эллиптико-параболалық тендеулер үшін орнатылған [3]. Көп-өлшемді эллиптико-параболалық тендеулер үшін жұмыс істеу түрі мен ретті өзгертумен рұқсат етілген және Дирихле есебін классикалық шешудің айқын түрі алынған.

**Түйін сөздер:** шешімділігі, аралас есеп, көп өлшемді эллиптико-параболалық тендеулер, Бессель функциясы.

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#### ЗАДАЧА ДИРИХЛЕ В ЦИЛИНДРИЧЕСКОЙ ОБЛАСТИ ДЛЯ ОДНОГО КЛАССА МНОГОМЕРНЫХ ЭЛЛИПТИКО-ПАРАБОЛИЧЕСКИХ УРАВНЕНИЙ

**Аннотация.** Краевые задачи для вырождающихся эллиптико-параболических уравнений на плоскости достаточно хорошо изучены ([1]). Корректность задачи Дирихле для вырожденных многомерных эллиптико-параболических уравнений с вырождением типа и порядка была установлена в [3]. В работе для многомерных эллиптико-параболических уравнений с вырождением типа и порядка показана разрешимость и получен явный вид классического решения задачи Дирихле.

**Ключевые слова:** разрешимость, смешанная задача, многомерные эллиптико-параболические уравнения, функция Бесселя.

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